## Optical Parametric Generation of Mid-Infrared Picosecond Pulses Beyond 6 µm in CdSiP<sub>2</sub>

O. Chalus, A. Esteban-Martin, P. G. Schunemann, K. T. Zawilski, J. Biegert, And M. Ebrahim-Zadeh,

<sup>1</sup>ICFO-Institut de Ciencies Fotoniques, Mediterranean Technology Park, 08860 Castelldefels, Barcelona, Spain

<sup>2</sup>BAE Systems, Inc., MER15-1813, P.O. Box 868, Nashua, New Hampshire 03061-0868, USA

<sup>3</sup>Institucio Catalana de Recerca i Estudis Avancats (ICREA), Passeig Lluis Companys 23, Barcelona 08010, Spain

<sup>4</sup>Department of Physics and Astronomy, University of New Mexico, 800 Yale Blvd, Albuquerque, New Mexico 87131-0001, USA

e-mail: majid.ebrahim@icfo.es

**Abstract:** We report parametric generation of near- and mid-infrared picosecond pulses at  $100 \, \text{kHz}$  in CdSiP<sub>2</sub> pumped at  $1.064 \, \mu \text{m}$ , providing  $154 \, \text{mW}$  of idler at  $6.204 \, \mu \text{m}$  and  $1.16 \, \text{W}$  of signal at  $1.282 \, \mu \text{m}$ .

© 2011 Optical Society of America

OCIS codes: 190.4400 Nonlinear optics, materials; 190.4970 Parametric oscillators and amplifiers

In the absence of widely available solid-state laser sources in the mid-infrared (mid-IR), optical parametric down-conversion has been established as an effective technique for the generation of coherent radiation in this spectral range. Using oxide-based birefringent materials such as LiNbO<sub>3</sub>, KTiOAsO<sub>4</sub> and RbTiOAsO<sub>4</sub>, and periodically-poled crystals such as PPLN and PPRTA, spectral regions up to ~5 µm can be accessed, but the onset of multiphonon absorption sets a practical upper limit of ~4 µm for wavelength generation in such materials. CdSiP<sub>2</sub> (CSP) [1], is a recently discovered nonlinear material which offers unique properties for parametric down-conversion into the mid-IR. It offers a transparency above ~6.5 µm and a noncritical phase-matching (NCPM) capability with an effective nonlinear coefficient as high as  $d_{\text{eff}} = d_{36} = 84.5 \text{ pm/V}$  [2]. Importantly, CSP has a bandgap well below 1 µm, which permits pumping at 1.064 µm, and under type I ( $e \rightarrow oo$ ) parametric generation with NCPM can provide an idler wavelength beyond 6.4 µm, a spectral range of great interest for medical applications. In earlier studies, the potential of CSP for the generation of mid-IR radiation using direct pumping at 1.064 µm was demonstrated in pulsed nanosecond and pulsed picosecond optical parametric oscillators [4-6]. Here, we report efficient generation of picosecond pulses in near- and mid-IR in CSP at 100 kHz using single-pass parametric generation (OPG) pumped by a mode-locked Nd:YVO<sub>4</sub> laser at 1.064 µm [7]. We demonstrate an average signal power of 1.16 W at 1.282 µm and idler power of 154 mW at 6.204 µm for 6.1 W of pump power.

A schematic of the experimental setup is shown in Fig. 1. The pump source is a commercial mode-locked Nd:YVO<sub>4</sub> laser at  $1.0642~\mu m$ . It can deliver up to 40 W of average power at 100~kHz with a pulse energy of  $400~\mu J$ . The output beam has a diameter of 5 mm and the pulses have durations of 8.7 ps, with a spectral bandwidth of  $\sim 0.2$  nm. The beam has a quality factor  $M^2 \sim 1.1$  and the output power stability of < 0.5% RMS over 13 hours. Using a telescope, the pump beam is collimated to a  $\sim 500~\mu m$  diameter before the CSP crystal. The pump power and polarization are controlled using two half-wave plates and a polarizing beam-splitter cube. Pumping is single-pass.

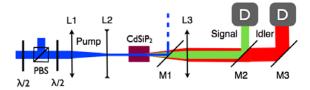


Fig. 1. Experimental setup for single-pass optical parametric generation in CSP pumped at 1.064 μm. λ/2: half-wave plate, PBS: polarizing beam-splitter, L: lens, M: mirror, D: diagnostics.

The CSP crystal was grown from a stoichiometric melt by the horizontal gradient freeze technique [8]. It was cut at  $\theta$ =90° ( $\varphi$ =45°) for type I ( $e\rightarrow oo$ ) interaction under NCPM with a length of 8 mm and an aperture of 6.75 mm x 6 mm (along the c-axis). The residual loss of the crystal was measured to be 0.198 cm<sup>-1</sup> for the pump at 1.064  $\mu$ m, 0.114 cm<sup>-1</sup> for the signal near 1.3  $\mu$ m, and 0.014 cm<sup>-1</sup> for the idler near 6.2  $\mu$ m. Both crystal faces were antireflection (AR)-coated for the three wavelengths with an eight-layer coating, providing an average reflectivity per surface of ~0.35% at 1064 nm, ~0.4% at 1.275  $\mu$ m, and ~0.5% at 6.2  $\mu$ m. The overall single-pass transmission of the AR-coated sample was 83% at 1.064  $\mu$ m.

## CTuD1.pdf

We observed OPG output at 1.1 W of pump, corresponding to a pulse energy of 11 μJ and pumping intensity of 0.62 GW/cm². By further increasing the pump to 6.1 W, we generated an average signal power of 1.16 W at 1.282 μm with pulse energy of 11.6 μJ. This represents a power efficiency of ~19% and a photon conversion of ~23%. At 6.1 W of pump, we measured an average idler power of 154 mW at 6.204 μm with a pulse energy of 1.54 μJ. Therefore, the power efficiency from the pump to the idler was ~2.5%, with a photon conversion efficiency as much as ~15%. Beyond 6.1 W of pump, we observed the onset of lensing in the CSP crystal. To ascertain the origin of this lensing effect, we chopped the pump beam and performed power scaling measurements of OPG output. The results are shown in Fig. 2, where the signal pulse energy is plotted against pump pulse energy, in the absence of chopping (I), and when the pump is chopped at 100 Hz with duty cycles of 50% (II) and 5% (III). The results clearly confirm the origin of the saturation and lensing as thermal. Given the high pump repetition rate of 100 kHz, saturation in output energy occurs as a result of heating of the crystal, leading to thermal dephasing at higher average powers. With crystals of higher quality, larger aperture and longer length, we expect substantial increases in the signal and idler power to multiwatt and watt level, and at higher conversion efficiencies, by increasing the available pump power to 40 W, while minimizing saturation and thermal lensing.

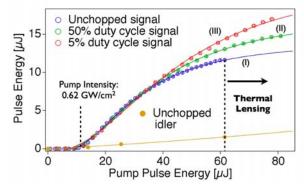


Fig. 2. Output signal pulse energy versus pump pulse energy at the input to the CSP crystal. (I): Unchopped pump beam; (II), (III): Chopped pump beam. Also shown is the idler pulse energy versus pump pulse energy.

The measured autocorrelation of the signal pulse and the corresponding spectrum are shown in Fig. 3(a) and 3(b), corresponding a signal pulse duration of 6.36 ps, and a spectrum centered at 1.282  $\mu$ m with a bandwidth of 8.5 nm, resulting in a time-bandwidth product of 9.3. The idler spectrum, shown in Fig. 3(c) centered at 6.204  $\mu$ m, and has a bandwidth of 122 nm. The dips in the spectrum correspond to absorption lines of water, as verified by the HITRAN molecular database. The signal and idler peak wavelengths of 1.282  $\mu$ m and 6.204  $\mu$ m are in close agreement with the calculated values of 1.286  $\mu$ m and 6.180  $\mu$ m for a pump wavelength of 1.0642  $\mu$ m based on the Sellmeier equations for the material [2].

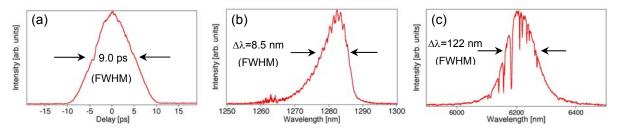


Fig. 3. (a) Intensity autocorrelation, and (b) spectrum of the signal pulses at 1.282 µm. (c) Spectrum of idler centered at 6.204 µm.

## References

- [1] P. G. Schunemann, K. T. Zawilski, T. M. Pollak, D. E. Zelmon, N. C. Fernelius, and F. Kenneth Hopkins, in *Advanced Solid-State Photonics*, Conference Program and Technical Digest (Optical Society of America, 2008), postdeadline paper MG6.
- [2] P. G. Schunemann, K. T. Zawilski, T. M. Pollak, V. Petrov, and D. E. Zelmon, in *Advanced Solid-State Photonics*, Conference Program and Technical Digest (Optical Society of America, 2009), paper TuC6.
- [4] V. Petrov, P. G. Schunemann, K. T. Zawilski, and T. M. Pollak, Opt. Lett. 34, 2399 (2009).
- [5] A. Peremans, D. Lis, F. Cecchet, P. G. Schunemann, K. T. Zawilski, and V. Petrov, Opt. Lett. 34, 3053 (2009).
- [6] V. Petrov, G. Marchev, P. G. Schunemann, A. Tyazhev, K. T. Zawilski, and T. M. Pollak, Opt. Lett. 35, 1230 (2010).
- [7] O. Chalus, P. G. Schunemann, K. T. Zawilski, J. Biegert, and M. Ebrahim-Zadeh, Opt. Lett. 35 (2010).
- [8] K. T. Zawilski, Peter G. Schunemann, T. M. Pollak, D. E. Zelmon, N. C. Fernelius, and F. K. Hopkins, J. Crystal Growth 312, 1127 (2010).